## Large magnetic entropy change and low hysteresis loss in the Nd- and Co-doped La(Fe,Si)<sub>13</sub> compounds

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The effect of Nd and Co substitution on magnetic entropy changes and hysteresis losses has been investigated for the cubic NaZn<sub>13</sub>-type LaFe<sub>13-x</sub>Si<sub>x</sub> compounds. Partially replacing La with Nd leads to a decrease of the Curie temperature  $T_C$  and an increase of the magnetic entropy change  $\Delta S$ . Substitution of Co for Fe in La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.5</sub>Si<sub>1.5</sub> can adjust  $T_C$  to around room temperature. A large  $\Delta S$  of 15 J/Kg K at  $T_C$ =280 K for a field change from 0 to 5 T and a small hysteresis loss close to zero near  $T_C$  have been obtained in La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub>. The Co-doped NaZn<sub>13</sub>-type LaNdFeSi compounds may be a suitable candidate for magnetic refrigerant near room temperature. © 2008 American Institute of Physics. [DOI: 10.1063/1.2829035]

Recently, many investigations have been carried out on the magnetocaloric effect (MCE) of the  $LaFe_{13-x}Si_x$  based compounds because of their potential application in magnetic refrigeration.<sup>1,2</sup> It was found that the LaFe<sub>13-x</sub>Si<sub>x</sub> compounds with a low Si concentration show a large magnetic entropy change due to the itinerant electron metamagnetic (IEM) transition.<sup>3-6</sup> Although the MCE of  $LaFe_{13-x}Si_x$  enhances with the decrease of Si content x, the Curie temperature  $T_C$ reduces, and it is usually much lower than the room temperature.<sup>7,8</sup> In order to work as a magnetic refrigerant near the ambient temperature, it is needed to adjust  $T_C$  to room temperature while retaining its large magnetic entropy change. Magnetic hysteresis loss is inevitable for the materials experiencing a first-order transition. To improve the efficiency of magnetic refrigeration, it is necessary to depress the hysteresis loss. A recent study has indicated that partially replacing La with Ce, Pr, and Nd in the LaFe<sub>13-r</sub>Si<sub>r</sub> compounds can enhance remarkably the MCE effect,<sup>9–11</sup> but also leads to a large hysteresis loss. In this paper, we report the effect of Nd and Co substitution on the magnetic entropy change  $\Delta S$  and hysteresis loss in the La(Fe, Si)<sub>13</sub> compounds. A large magnetic entropy change and a low hysteresis loss can be obtained in the Co-doped NaZn<sub>13</sub>-type LaNdFeSi compounds.

 $\begin{array}{cccc} Samples & of & LaFe_{11.5}Si_{1.5}, & LaFe_{11.2}Si_{1.8}, \\ La_{0.7}Nd_{0.3}Fe_{11.5}Si_{1.5}, & La_{0.7}Nd_{0.3}Fe_{11.2}Si_{1.8}, & \text{and} \\ La_{0.7}Nd_{0.3}Fe_{10.7}Co_{0.8}Si_{1.5} & \text{were prepared by arc melting in an} \\ argon atmosphere of high purity. The purity of starting ma-$ 

terials is 99% for Nd, 99.9% for La, Fe, and Co, and 99.999% for Si. The as-prepared ingots were wrapped by molybdenum foil, sealed in a quartz tube of high vacuum, annealed at 1373 K for 40 days, and then quenched to room temperature. X-ray diffraction (XRD) measurements on powder samples were performed using Cu  $K\alpha$  radiation to identify the phase structure and the crystal lattice parameter. Magnetizations were measured as functions of temperature and magnetic field by using a superconducting quantum interference device magnetometer. The isothermal magnetic entropy change was calculated from the magnetization data by using the Maxwell relation.

Figure 1 shows the room-temperature powder XRD patterns of LaFe<sub>11.5</sub>Si<sub>1.5</sub>, LaFe<sub>11.2</sub>Si<sub>1.8</sub>, La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub>, La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.2</sub>Si<sub>1.8</sub>, and La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub>. These compounds crystallized in a very clean single phase of a cubic NaZn<sub>13</sub>-type structure. The lattice parameter *a* ob-

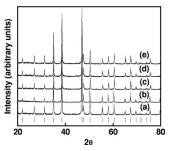


FIG. 1. Room-temperature powder XRD patterns of LaFe<sub>11.5</sub>Si<sub>1.5</sub> (a), LaFe<sub>11.2</sub>Si<sub>1.8</sub> (b), La<sub>07</sub>Nd<sub>0.3</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub> (c), La<sub>07</sub>Nd<sub>0.3</sub>Fe<sub>11.2</sub>Si<sub>1.8</sub> (d), and La<sub>07</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub> (e). Bragg reflections (small vertical lines) are also shown.

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TABLE I. The lattice parameter *a*, Curie temperature  $T_c$ , and magnetic entropy change  $\Delta S$  for the magnetic field changes of 0–2 and 0–5 T, and the magnetic hysteresis losses for LaFe<sub>11.5</sub>Si<sub>1.5</sub>, LaFe<sub>11.2</sub>Si<sub>1.8</sub>, La<sub>07</sub>Nd<sub>0.3</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub>, La<sub>07</sub>Nd<sub>0.3</sub>Fe<sub>11.2</sub>Si<sub>1.8</sub>, and La<sub>07</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub>.

Compounds	a (Å)	<i>Т<sub>С</sub></i> (К)	ΔS (0–2 T) (J/kg K)	ΔS (0–5 T) (J/kg K)	Hysteresis loss (J/kg)
LaFe <sub>11.5</sub> Si <sub>1.5</sub>	11.4686	194	20.9	23.7	21.2
LaFe <sub>11.2</sub> Si <sub>1.8</sub>	11.4635	216	7.8	13.7	No
La <sub>07</sub> Nd <sub>0.3</sub> Fe <sub>11.5</sub> Si <sub>1.5</sub>	11.4502	188	29.3	32.0	78.1
La07Nd0.3Fe11.2Si1.8	11.4426	207	10.5	15.2	3.5
La <sub>0.7</sub> Nd <sub>0.3</sub> Fe <sub>10.7</sub> Co <sub>0.8</sub> Si <sub>1.5</sub>	11.4533	280	7.9	15.0	No

tained from the XRD patterns is listed in Table I. It is found that the substitution of Nd leads to a contraction of the lattice, as has been observed in  $La_{1-x}Ce_xFe_{13-y}Si_y$  compounds.<sup>9</sup>

Figure 2 shows the thermomagnetic M-T curves for LaFe<sub>11.5</sub>Si<sub>1.5</sub>, La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub>, LaFe<sub>11.2</sub>Si<sub>1.8</sub>,  $La_{0.7}Nd_{0.3}Fe_{11.2}Si_{1.8}$ , and  $La_{0.7}Nd_{0.3}Fe_{10.7}Co_{0.8}Si_{1.5}$  compounds measured in the heating and cooling processes under a magnetic field of 0.01 T. For LaFe<sub>11.5</sub>Si<sub>1.5</sub> and La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub>, an obvious temperature hysteresis is observed between the transition on heating and cooling, indicating that two samples have a thermal-induced first-order magnetic transition at  $T_C$ . A nearly reversible change of magnetization with temperature is observed for LaFe<sub>11.2</sub>Si<sub>1.8</sub>, La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.2</sub>Si<sub>1.8</sub>, and La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub>, which is a characteristic of second-order magnetic transition or a weak first-order magnetic transition. It can be seen from Fig. 2 that the substitution of Si and/or Co for Fe drives the magnetic transition from first order to second order. The Curie temperature  $T_C$  is determined from the thermomagnetic M-T curves obtained in an external magnetic field H =0.01 T. Table I summarizes the values of  $T_C$ . One can find that substitution of Nd for La downward shifts  $T_C$ , while the substitution of Co for Fe drives  $T_C$  upward. The small decrease of  $T_C$  with increasing Nd content is due to the lattice contraction, as observed in the Ce-doped  $LaFe_{13-x}Si_x$ compounds.9 A compound La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub> with  $T_C$ =280 K is obtained. The obvious enhancement of  $T_C$  may result from the contributions of the strong Fe-Co interactions. Our study shows that in La<sub>1-x</sub>Nd<sub>x</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub> compounds, properly replacing Fe with Co can adjust  $T_C$  to around room temperature, which is essential for roomtemperature magnetic refrigeration application.

As an example, Fig. 3 shows the isothermal magnetiza-

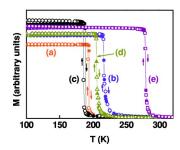


FIG. 2. (Color online) Temperature dependence of the magnetization measured on heating and cooling in a magnetic field of 0.01 T for LaFe<sub>11.5</sub>Si<sub>1.5</sub> (a), LaFe<sub>11.2</sub>Si<sub>1.8</sub> (b), La<sub>07</sub>Nd<sub>0.3</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub> (c), La<sub>07</sub>Nd<sub>0.3</sub>Fe<sub>11.2</sub>Si<sub>1.8</sub> (d), and La<sub>07</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub> (e).

tion curves for La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub> and La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub> in a wide temperature range around the Curie temperature with different temperature steps in magnetic fields up to 5.0 T. The temperature step of 2 K is chosen in the vicinity of  $T_C$  and steps of 5 K for the regions far away from  $T_C$ . It can be seen from Fig. 3(a) that an obvious magnetic hystersis loop is observed for La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub>, indicating a characteristic of the IEM transition and the nature of a first-order magnetic transition, as found in LaFe<sub>11.7</sub>Si<sub>1.3</sub>.<sup>3-6,12</sup> For La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub>, however, no magnetic hystersis is observed, as shown in Fig. 3(b); that is, the Co-doped NaZn<sub>13</sub>-type LaNdFeSi compounds show a reversible change of the magnetization with temperature and magnetic field. This is very favorable to the magnetic refrigeration application since a completely reversible MCE requires that there exists no hysteresis as the magnetization varies with temperature and magnetic field.

Figure 4 shows the temperature dependences of hysteresis loss for  $La_{0.7}Nd_{0.3}Fe_{11.5}Si_{1.5}$  and  $La_{0.7}Nd_{0.3}Fe_{10.7}Co_{0.8}Si_{1.5}$ , respectively. The hysteretic loss is defined as the enclosed area between the ascending and descending branches of magnetization curve. Table I lists the maximum values of hysteresis loss for  $LaFe_{11.5}Si_{1.5}$ ,

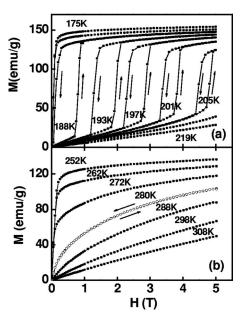


FIG. 3. Isothermal magnetization curves of  $La_{07}Nd_{0.3}Fe_{11.5}Si_{1.5}$  (a) and  $La_{07}Nd_{0.3}Fe_{10.7}Co_{0.8}Si_{1.5}$  (b) on field increase and decrease around the Curie temperature.

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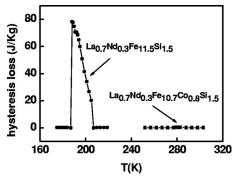


FIG. 4. Temperature dependence of the hysteresis loss of  $La_{07}Nd_{0.3}Fe_{11.5}Si_{1.5}$  and  $La_{07}Nd_{0.3}Fe_{10.7}Co_{0.8}Si_{1.5}$ .

LaFe<sub>11.2</sub>Si<sub>1.8</sub>, La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub>, La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.2</sub>Si<sub>1.8</sub>, and La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub>. For La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.5</sub>Si<sub>1.5</sub>, a large hysteresis loss of 78.1 J/kg is observed, while the La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub> exhibits a hysteresis loss close to zero. It is clear that the magnetic first-order phase transition is suppressed by the substitution of Co for Fe, leading to the disappearance of hysteresis losses.

The magnetic entropy change  $\Delta S$  is calculated from magnetization data by using the following equation:

$$\Delta S(T,H) = \int_0^H \left(\frac{\partial M}{\partial T}\right)_H dH,$$

which is based on the Maxwell relation. The  $\Delta S$  for La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.2</sub>Si<sub>1.8</sub> and La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub> compounds as a function of temperature for different magnetic field changes are shown in Fig. 5. Table I lists the maximum values of  $\Delta S$  for the magnetic field changes from 0 to 2 T and 0 to 5 T at  $T_C$ . It is found that substitution of Nd for La in LaFe<sub>13-x</sub>Si<sub>x</sub> leads to a remarkable increase of magnetic entropy change. The enhancement of  $\Delta S$  is attributed to the strengthening of IEM transition above  $T_C$  caused by the substitution of Nd, similar to the case of Ce-doped LaFeSi.<sup>9</sup> Figures 6(a) and 6(b) show the Arrott plots of La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.2</sub>Si<sub>1.8</sub> and La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub> compounds, respectively. An obvious inflection point in the Arrott plots at  $T_C$  for La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>11.2</sub>Si<sub>1.8</sub> is the signature of the IEM transition from paramagnetic to ferromagnetic order

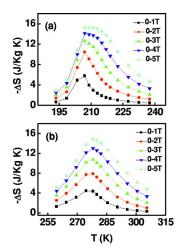


FIG. 5. (Color online) Temperature dependence of the magnetic entropy change of  $La_{07}Nd_{0.3}Fe_{11.2}Si_{1.8}$  (a) and  $La_{07}Nd_{0.3}Fe_{10.7}Co_{0.8}Si_{1.5}$  (b) for the magnetic field changes of 0–1, 0–2, 0–3, 0–4, and 0–5 T.

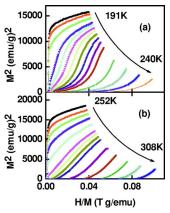


FIG. 6. (Color online) Arrott plots of  $La_{0.7}Nd_{0.3}Fe_{11.2}Si_{1.8}$  (a) and  $La_{0.7}Nd_{0.3}Fe_{10.7}Co_{0.8}Si_{1.5}$  (b).

above  $T_C$ . For La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub>, a characteristic of second-order transition is observed, as shown in Fig. 6(b). The maximum value of  $\Delta S$  for La<sub>0.7</sub>Nd<sub>0.3</sub>Fe<sub>10.7</sub>Co<sub>0.8</sub>Si<sub>1.5</sub> is 15 J/kg K at  $T_C$  for a field change of 0–5 T, which is larger than that of Gd.<sup>13</sup> Since the highest magnetocaloric effect involving a second-order magnetic transition near room temperature is produced by Gd, the Co-doped La<sub>1-r</sub>Nd<sub>r</sub>Fe<sub>115</sub>Si<sub>15</sub> compounds are attractive candidates for magnetic refrigerants in an extended high temperature range even at room temperature.

We have studied the magnetic entropy change  $\Delta S$  and the magnetic hysteresis loss in the Nd- and Co-doped cubic  $NaZn_{13}$ -type La(Fe, Si)<sub>13</sub> compounds. It is found that substitution of Nd for La downward shifts  $T_C$ , while the substitution of Co for Fe upward shifts  $T_C$ . The substitution of Nd for La enhances the characteristic of the IEM transition in a large  $\Delta S$ . above  $T_C$ , resulting For  $La_{0.7}Nd_{0.3}Fe_{10.7}Co_{0.8}Si_{1.5}$ , the maximum value of  $\Delta S$  is 15 J/kg K at  $T_C$ =280 K for a field change of 0–5 T, which is larger than that of Gd. A remarkable result is that the magnetic behavior of the Co-doped compounds is nearly reversible for the field and temperature increase-decrease cycling, and no hysteresis loss near  $T_C$  is observed. We believe that the Co-doped NaZn<sub>13</sub>-type LaNdFeSi compounds may be a suitable candidate for magnetic refrigerant near room temperature.

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- <sup>1</sup>K. A. Gschneidner, Jr. et al., Rep. Prog. Phys. 68, 1479 (2005).
- <sup>2</sup>A. M. Tishin and Y. I. Spichkin, in *The Magnetocaloric Effect and Its Applications*, edited J. M. D. Coey *et al.* (Institute of Physics, University of Reading, Berkshire, 2003).
- <sup>3</sup>F. X. Hu et al., Chin. Phys. 9, 550 (2000).
- <sup>4</sup>F. X. Hu *et al.*, Appl. Phys. Lett. **78**, 3675 (2001).
- <sup>5</sup>S. Fujieda et al., Appl. Phys. Lett. 81, 1276 (2002).
- <sup>6</sup>S. Fujieda et al., Sci. Technol. Adv. Mater. 4, 339 (2003).
- <sup>7</sup>T. T. M. Palstra et al., J. Magn. Magn. Mater. 36, 290 (1983).
- <sup>8</sup>A. Fujita, Y. Akamatsu, and K. Fukamichi, J. Appl. Phys. 85, 4756 (1999).
- <sup>9</sup>S. Fujieda et al., J. Alloys Compd. **408**, 1165 (2006).
- <sup>10</sup>S. Fujieda et al., J. Magn. Magn. Mater. **310**, e1004 (2007).
- <sup>11</sup>D. T. Kim Anh et al., J. Magn. Magn. Mater. 262, 427 (2003).
- <sup>12</sup>N. L. Di et al., Phys. Rev. B 69, 224411 (2004).
- <sup>13</sup>S. Yu Dan'kov et al., Phys. Rev. B 57, 3478 (1998).